## BRIEF COMMUNICATIONS

## TEMPERATURE DISTRIBUTION IN A LIQUID LAYER

ON A HORIZONTAL SOLID SURFACE
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The temperature distribution is found in an incompressible still liquid of known mass covering a horizontal solid surface at constant temperature. The exact layer thickness and the total heat content are determined. A stability condition is derived for the liquid equilibrium.

We consider a liquid layer of known mass on a solid horizontal surface whose temperature, like that of the external medium at the free surface of the liquid, is constant and also known. For a still liquid, we have the system of equations

$$
\begin{equation*}
\frac{d}{d z}\left(\lambda \frac{d t}{d z}\right)=0, \quad \int_{0}^{h} \frac{d z}{v}=m \tag{1}
\end{equation*}
$$

At constant thermal conductivity, the liquid temperature is a linear function of the coordinate $t=t_{1}+a z$. Its parameters and $h$ are to be determined.

Analogous problems have been discussed for a liquid layer covering a gravitating sphere [1] and on the inside of a rotating cylinder [2].

If the $\alpha_{i}$ are assumed constant, the boundary conditions on the temperature yield

$$
\begin{equation*}
a=\gamma\left(t_{1}-t_{e 1}\right), \quad t_{1}+\beta x=\theta \tag{2}
\end{equation*}
$$

In particular, we have $t=t_{e_{2}}$ when $\alpha_{1}=0, t=t_{e 1}$ when $\alpha_{2}=0$, and $t=t_{e}$ when $\Delta t_{e}=0$; we assume below that these cases are excluded. Corresponding to these cases we have

$$
\min \left(t_{e 1}, t_{e 2}\right)<t, \theta<\max \left(t_{e 1}^{\prime}, t_{e 2}\right), \operatorname{sgn} a=\operatorname{sgn} \Delta t_{e}
$$

We restrict the discussion to an incompressible liquid with a constant coefficient of thermal expansion:

$$
\begin{equation*}
v=v_{0} \exp (\delta t), \quad v_{0}=\text { const. } \tag{3}
\end{equation*}
$$

After $t_{1}$ and $a$ are eliminated with the help of (2), the integral relation in (1) becomes

$$
\begin{equation*}
[1-\exp (-u)] \exp (\beta u)=b(\varepsilon-u) \tag{4}
\end{equation*}
$$

When $\alpha_{2}=\alpha_{1}$, we have

$$
\begin{equation*}
\operatorname{sh} \frac{u}{2}=2 b(\varepsilon-u) ; 2 b=\gamma h_{0} \exp \left(\frac{t_{e 1}+t_{e 2}}{2}\right) . \tag{5}
\end{equation*}
$$

At $\delta=0$, we easily find the final solution:

$$
\begin{equation*}
t_{1}=t_{e 1}+\frac{\beta}{1+b} \Delta t_{e}, \quad a=\frac{\beta \gamma}{1+b} \Delta t_{e}, \quad h=h_{0} \quad\left(b=\beta \gamma h_{0}\right) . \tag{6}
\end{equation*}
$$

To find $x$, we must in general solve transcendental equation (4), which has a single root in the range $0<x / \Delta t_{e}<1$. All the unknown quantities are found from $x$ :

$$
\begin{equation*}
t_{1}=t_{e 1}+\beta y, \quad a=\beta \gamma y, \quad h=x /(\beta \gamma y) . \tag{7}
\end{equation*}
$$

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The assumed equilibrium of the still liquid is stable when $a>-\mathrm{g} \delta \mathrm{t} / \mathrm{c}$ [3]; i.e., it is always stable if $\Delta \mathrm{t}_{\mathrm{e}}>0$, or it is stable when $|a|<\mathrm{g} \delta \mathrm{t}_{2} / \mathrm{c}$ if $\Delta \mathrm{t}_{\mathrm{e}}<0$. Otherwise, convection may occur, generally causing a temperature distribution different from that obtained.

The total heat content of a liquid column with a base of unit area is

$$
\begin{equation*}
i=c \int_{0}^{h} \frac{t}{v} d z=\frac{c m}{\delta}\left(1-\beta u+\frac{\beta^{\prime}-\beta}{\exp u-1}\right) ; \tag{8}
\end{equation*}
$$

when $\delta=0$, we have

$$
\begin{equation*}
i=c m\left[t_{e 1}+(\beta+b / 2) \frac{\Delta t_{e}}{1+b}\right] \tag{9}
\end{equation*}
$$

Let us consider the case of small $\delta(|\varepsilon| \ll 1)$ separately. Expanding the exponents in (4) in series and discarding powers of $u(|u|<|\varepsilon|)$ greater than the second, we find the quadratic equation

$$
\begin{equation*}
(\beta-1 / 2) u^{2}+(1-b) u-b \varepsilon=0 . \tag{10}
\end{equation*}
$$

Its roots are real. The root which vanishes along with $\Delta t_{\mathrm{e}}$ is

$$
\begin{equation*}
x=\left[1+\frac{\beta^{\prime}-\beta}{2} \frac{b \delta}{(1+b)^{2}} \Delta t_{z}\right] \frac{b}{1+b} \Delta t_{e} . \tag{11}
\end{equation*}
$$

We find from (7) that

$$
\begin{gather*}
t_{1}=t_{\varepsilon 1}+\left[1+\frac{\beta^{\prime}-\beta}{2}\left(\frac{b}{1+b}\right)^{2} \varepsilon\right] \frac{\beta}{1+b} \Delta t_{e}, \\
a=\left[1+\frac{\beta-\beta^{\prime}}{2}\left(\frac{b}{1+b}\right)^{2} \varepsilon\right] \frac{\beta \gamma}{1+b} \Delta t_{e},  \tag{12}\\
h=h_{0} \exp (\delta \theta)\left(1+\frac{\beta^{\prime}-\beta}{2} \frac{b}{1+b} \varepsilon\right) . \tag{13}
\end{gather*}
$$

When $\delta=0$, Eqs . (12) and (13) convert into Eqs . (6).

## NOTATION

$\mathrm{v}, \mathrm{c}, \lambda, \delta \quad$ are the specific volume, heat capacity, thermal conductivity, and coefficient of thermal expansion, respectively, of the liquid;
$\alpha \quad$ is the heat-exchange coefficient (the contact thermal conductivity [1, 2, 4]);
$\mathrm{m}, \mathrm{i}$ are the mass and total heat content of a liquid column with a base of unit area;
$\mathrm{Z} \quad$ is the coordinate, equal to 0 at the solid surface and $h$ at the free surface of the liquid;
$\mathrm{t}, a$ are the absolute temperature and temperature gradient, respectively;
$\gamma=\alpha_{1} / \lambda$;
$\beta=1-\beta^{\prime}=\alpha_{2} /\left(\alpha_{1}+\alpha_{2}\right) ;$
$\mathrm{h}_{0}=\mathrm{mv}_{0}$;
$\Delta t_{e}=t_{e_{2}}-t_{e_{1}}$;
$\varepsilon=\delta \Delta t_{\mathrm{e}}$;
$\mathrm{x}=\mathrm{t}_{2}-\mathrm{t}_{1}=a \mathrm{~h}$;
$\mathrm{y}=\Delta \mathrm{t}_{\mathrm{e}}-\mathrm{x}$;
$\mathrm{u}=\delta \mathrm{x}$;
$\theta=\beta^{\prime} \mathrm{t}_{\mathrm{e}_{1}}+\beta \mathrm{t}_{\mathrm{e}_{2}} ;$
$\mathrm{b}=\beta \gamma \mathrm{h}_{0} \exp (\delta \theta)$.

## Subscripts

e, 1, 2 are the external medium, solid surface, and free surface of the liquid, respectively.

## LITERATURE CITED

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